

Terahertz to UV Generalized Magneto-optic ellipsometry on ZnMnSe: Giant Kerr effect, band-to-band transitions, and charge transport parameters



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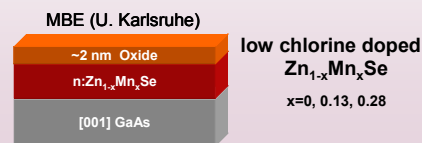
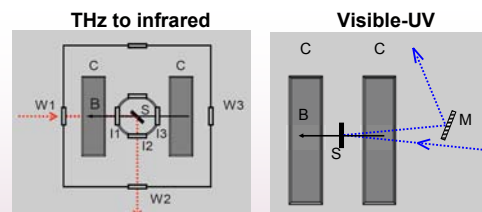
Our message

- I. **Kerr rotation in paramagnetic ZnMnSe.**
- II. **Optical anisotropy found in ZnMnSe, probably due to in-plane strain, or atomic ordering.**
- III. **Zeeman splitting of the conduction bands as function of Mn content.**
- IV. **Free charge carrier parameters of ZnMnSe.**

Experiment and model

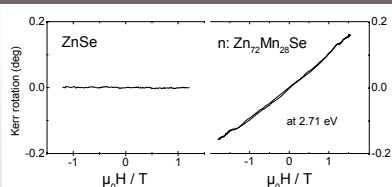
- ✓ Magneto-Optic Generalized Ellipsometry (MOGE) at RT on ZnSe and Zn_{1-x}Mn_xSe (x=0, 0.13, 0.28) in the spectral range from 70 to 650 cm⁻¹ and from 1.2 to 3.4 eV.
- ✓ Isotropic model dielectric function (MDF) for ZnSe, and ZnMnSe (Phys. Rev. B **70**, 04513, (2004))
- ✓ Observation of in-plane anisotropy in ZnMnSe (in absence of a magnetic field)
- ✓ Anisotropic MDF of ZnMnSe
- ✓ MOGE (Kerr rotation) measurements of ZnSe and Zn_{1-x}Mn_xSe (x=0, 0.13, 0.28)
- ✓ MDF for Zeeman splitting in ZnMnSe (magnetic field induced birefringence)
- ✓ Charge transport parameters of low chlorine doped ZnMnSe

Setup



Experimental results

Kerr rotation measurements



MOGE measurements in quasi-Kerr configuration for ZnSe, and Zn₇₂Mn₂₈Se. The Zn₇₂Mn₂₈Se sample shows a paramagnetic response for a photon energy corresponding to the band-to-band transition, while in ZnSe no Kerr rotation is detected.

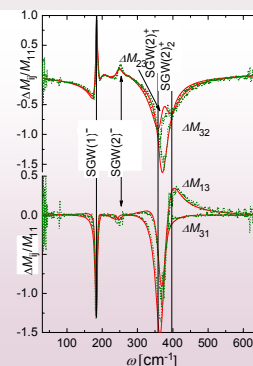
Free charge carrier parameters

Γ-Point CB effective mass

	N [10 ¹⁷ cm ⁻³]	m^* [m_0]	μ [10 ² cm ² /Vs]
GaAs-sub.	10.5(1)	0.071(1)	20.5(1)
ZnMnSe	4.9(2)	0.086(2)	3.0(2)

corresponds to kp-calculations of the Γ-Point CB effective mass for Zn_{0.87}Mn_{0.13}Se

Hofmann, Schade, Schubert, et al., Appl. Phys. Lett. **88**, 042105 (2006)



Isotropic

In-plane birefringence

Magnetic induced birefringence

$$\epsilon(\omega) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \frac{1}{2}(\epsilon_{\parallel} + \epsilon_{\perp}) + \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \frac{1}{2}(\epsilon_{\parallel} - \epsilon_{\perp}) + \begin{pmatrix} 0 & i & 0 \\ -i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \frac{1}{2}(\epsilon_{+} - \epsilon_{-})$$

Dielectric response

M_{12}

$\frac{1}{2}(\epsilon_{\parallel} + \epsilon_{\perp})$

M_{14}

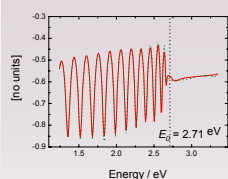
$\frac{1}{2}(\epsilon_{\parallel} - \epsilon_{\perp})$

M_{12}

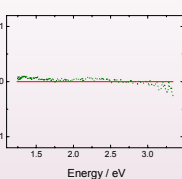
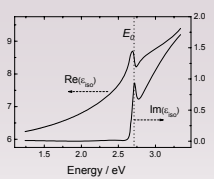
$\frac{1}{2}(\epsilon_{+} - \epsilon_{-})$

Materials

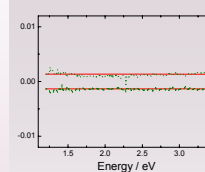
ZnSe



The real and imaginary parts of the isotropic dielectric function obtained from the best match analysis of the GE experimental data.

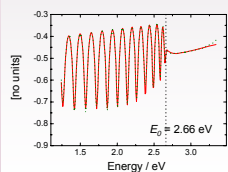


The Mueller matrix element M_{14} used to detect in-plane birefringence shows no signal in the studied spectral range.

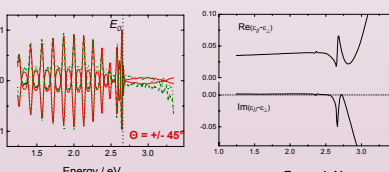
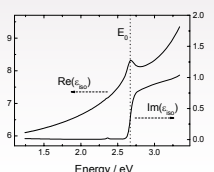


No signal was detectable from the MOGE experiment as shown by the M_{23} and M_{32} elements.

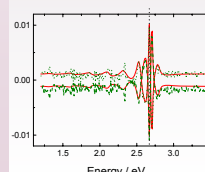
Zn72Mn28Se



Same as above.



A significant difference between the parallel and perpendicular dielectric function was found for ZnMnSe.



12 meV conduction band splitting energy detected in ϵ_{\parallel} and ϵ_{\perp} at $\mu_0 H = 1.8$ T.